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# MIXED ALKALI EFFECT ON STRUCTURE AND OPTICAL PROPERTIES OF QUATERNARY B<sub>2</sub>O<sub>3</sub>-P<sub>2</sub>O<sub>5</sub>-LI<sub>2</sub>O-CAO GLASS CONTAINING FE<sub>2</sub>O<sub>3</sub>

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### **ABSTRACT**

Series of alkali borophosphate glasses containing  $Fe_2O_3$  is prepared by the conventional melt-quenching technique. The alkali mixed effect on the glass structure and its optical properties is studied through replacing  $Li_2O$  (alkali oxide) by CaO (alkaline-earth oxide). Positron spectroscopy revealed a periodic change in both the glass refractive indices and free volume size as the molar ratio of  $Li_2O$  increases. It is an evidence for a variation of nonbridging oxygen bonds (NBO) density. The molar polarizability of NBO explains the obtained optical properties, while the variation in the glass molar volume clarifies the detected variation in the glass free volume. FT-IR spectroscopy of the investigated glasses confirmed the variation in the relative NBO density as  $Li_2O$  content increases on the expense of CaO.

**KEYWORDS:** Mixed Alkali Effect on Structure and Optical Properties, Alkali Borophosphate, Fe<sub>2</sub>O<sub>3</sub>, Li<sub>2</sub>O, CaO

#### INTRODUCTION

Due their low melting temperature, high thermal expansion, wide ultraviolet transparency and low dispersion compared with silicate glasses, phosphate glasses are of great interest for various optical applications such as optical fibers [1]. Employ of phosphate glasses as achromatic host materials in solid state lasers depends on types of dopants and their concentrations [2]. However, the use of phosphate glasses in solid state lasers was limited by their low chemical durability [3].

The basic building blocks of phosphates are the P-tetrahedra which link through covalent bridging oxygens to form various phosphate anions ( $\nu$ -P<sub>2</sub>O<sub>5</sub>). The addition of trivalent cations into metaphosphate (ratio of O/P = 3.0) glasses can lead to an increase in the degree of structure connectivity through the substitute of P=O nonbridging bonds by bridging bonds [4]. For that reason, B<sub>2</sub>O<sub>3</sub> is added in the form of BO<sub>4</sub> tetrahedra into phosphate glasses to improve their chemical durability providing more three-dimensional cross-linked structures [5-8]. On the other hand, borates are of exceptional importance due to their interesting linear and nonlinear optical properties [9-11]. Borophosphate glasses exhibit flat gain broadband amplification [12]. In addition, alkali-earthed-borophosphate glass matrices are of interest for many applications such as hosts for rare-earth dopants for fiber amplifiers [13]. Also, alkali-borophosphate glass exhibits large electro-optical Kerr-like effect and so large third-order nonlinearity [14, 15]. Addition of B<sub>2</sub>O<sub>3</sub> to a phosphate glass improves the chemical durability as well as thermal and mechanical stability of the pure phosphate glass [16-19]. Fe<sub>2</sub>O<sub>3</sub> has excellent effect on chemical durability in phosphate glass systems. The good chemical durability of the iron phosphate glass was attributed to a large number of Fe-O-P bonds in the structure [20]. Addition of iron to borate glass makes it electrically semiconducting [21]. Phosphate glasses containing divalent iron cations exhibit interesting semiconducting properties, magneto-optical properties, magnetic transition properties and electro-optic devices [22, 23].

As functions of  $Fe_2O_3$  content, iron ions exist in different valence states with different local symmetries in the glass matrices. For example, iron ions exist as  $Fe^{3+}$  with both tetrahedral and octahedral coordination, and/or as  $Fe^{2+}$  with octahedral coordination. Both of  $Fe^{3+}$  and  $Fe^{2+}$  are paramagnetic ions [24]. The coexistence of two valence states of iron leads to interesting semiconducting and magnetic properties for iron phosphate glasses [25].

Several glasses exhibit a linear-like behavior of their physical properties with varying chemical composition. An exception exists for glasses containing two different alkali oxides where large deviation from linearity occurs having maxima or minima. The mixed alkali effect (MAE) was observed for instant as nonlinear variation in electrical conductivity isotherms when a fraction of one type of the mobile ions is replaced by another type, keeping the total alkali content constant [26]. In this work the mixed alkali effect is used to study maxima for refractive index and minima in positron lifetime are studied with the prepared glasses.

#### **EXPERIMENTAL**

#### **Glass Preparation**

Series of borophosphate glass containing iron oxide of the composition  $(0.3B_2O_3 - 0.3P_2O_5 - 0.2Fe_2O_3 - xLi_2O - (0.2 - x) CaO)$  with  $0 \le x \le 0.2$  mol fraction is prepared. The used materials are of chemically pure grade, in the form of Fe<sub>3</sub>O<sub>4</sub>, H<sub>3</sub>BO<sub>3</sub>, CaO, 2[NH<sub>4</sub>H<sub>2</sub>PO<sub>4</sub>] and Li<sub>2</sub>CO<sub>3</sub>. The amount of the glass batch is 50 g melt<sup>-1</sup>. The glass is prepared by melt quenching technique using platinum 2% rhodium crucibles in an electric furnace. The batch was pre-heated at 500-600 °C for almost an hour to evaporate the carbonates. The temperature of melting is 1000-1100 °C, the duration of melting is one hour after the last traces of batches are disappeared.

To avoid the presence of bubbles the glasses are continue stirred during the glass melt preparation. Then the melt is poured onto stainless steel mould and annealed at around  $350^{\circ}$ C to remove the thermal strains. Optical slabs are prepared by grinding and polishing of the prepared samples with paraffin oil and minimum amount of water. The thickness of the glass slabs is about 3 mm. Polishing method is completed with stannic oxide and paraffin to reach a surface roughness less than  $\lambda/3$ , which is tested by interferometric method. The homogeneity of the glasses is examined using two crossed polarizers.

### **Spectrophotometric Measurements**

Computer aided two-beam spectrophotometer (shimadzu-3101PC UV-VIS NIR) is used to record the reflectance, R, and the transmittance, T, data of plane-parallel polymeric film samples, A resolution limit of 0.2 nm and a sampling interval of 2 nm are utilized for the different measuring points, The accuracy of measuring  $R(\lambda)$ , and  $T(\lambda)$  is 0.003 with the incident beam making an angle of  $5.0^{\circ}\pm0.1^{\circ}$  to the normal to external slab faces. The measurements are carried out at room temperature for the entire spectral range 0.2-2.5  $\mu$ m. The different optical applications of the glass require information about glass refractive index and its dispersion [9]. Applying Fresnel's theory, the glass refractive index, n, is determined by means of an iteration technique [27] using the transmittance and reflectance recorded data. The estimated uncertainty in the index of refraction is  $\pm$  0.005.

#### **Positron Annihilation Measurements**

Positron annihilation measurements are carried out at room temperature for the investigated glass samples. The PAL spectra are recorded using a fast-fast coincidence system with a time resolution of 250 picoseconds (ps). The

positron source is prepared by drying a droplet of <sup>22</sup>NaCl aqueous solution (about 20µCi) on a thin kapton foil. The positron source is sandwiched between two similar samples that are each about 1 mm thick; the two samples included in the positron source are put it between the two detectors of the positron lifetime set-up.

With the emission of a positron, the source emits a  $\gamma$ -quantum, a photon that is recorded on one detector. A  $\gamma$ - ray is emitted if an electron and its antiparticle, the positron, meet and annihilate, and one of the annihilation photons is recorded on another detector. The time difference between the two signals from the two detectors is, consequently, the positron lifetime. With a suitable electronics set-up, this time difference is transformed into what is called the lifetime spectra.

The lifetime spectrum shows a curve, which is a decaying exponential, containing many components due to several annihilation processes. The area under one component divided by the total area of the spectrum is called the relative intensity (I) of the component. A positron can annihilate through different processes, each of which gives rise to a certain mean lifetime ( $\tau$ ). The analyses of the positron annihilation lifetime spectra are performed using the PATFIT program [28], which gives the average of the lifetimes and their intensities. The positron lifetime technique is discussed elsewhere [28-30].

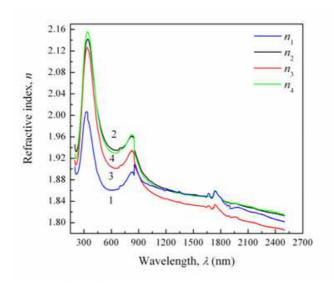
#### FT-IR Vibrational Absorption Measurements

Infrared vibrational absorption measurements are recorded for the present glass compositions in the range of (1400 - 400 cm<sup>-1</sup>) using Jasco FT/IR-300E infrared spectrophotometer. The alkali halide disc-technique at room temperature is used. Well-dried and ground glasses are mixed with well-dried infrared grade potassium bromide and then sufficiently ground to obtain a homogeneous mixture of minimum particle size. The mixture is mechanically pressed at 70-Mpa pressure in the form of discs.

#### RESULTS AND DISCUSSIONS

## **Dispersion Properties of Investigated Glasses**

Figure 1 illustrates the calculated dispersion of refractive indices in a spectral range from 0.2 to 2.5  $\mu$ m. An alternative change in the glass dispersion is seen as the ratio of Li<sub>2</sub>O/CaO increases. At percentage of relative molar ratio of Li<sub>2</sub>O/CaO = 25 % the glass refractive index increases  $n_d = 1.8604$ , as shown in Fig. 2. The refractive index,  $n_d$ , is measured at the standard wavelength  $\lambda = 587.6$  nm. At percentage of relative molar ratio of Li<sub>2</sub>O/CaO = 50 %, the glass refractive indices starts to increase having  $n_d = 1.9396$ . Moreover, decrease in glass refractive index is seen with sample 3 reaching  $n_d = 1.9056$  at percentage of relative molar ratio of Li<sub>2</sub>O/CaO = 75 %. With no contribution from CaO in the glass structure, i.e., at a percentage for the relative molar ratio of Li<sub>2</sub>O/CaO = 100 % the refractive index increased once again reaching  $n_d = 1.9335$ . Such alternation in the glass index is attributed to an fluctuation in oscillator strength value,  $f_{ii}$ .



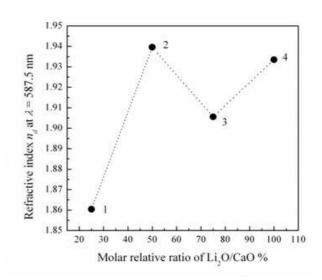


Figure 1: Dispersion of Glass Refractive Indices against Wavelength

Figure 2: Variation of Glass Refractive Indices at  $\lambda = 587.6$  nm versus the Percentage of Molar Ratio of Li<sub>2</sub>O/CaO

As it will be seen later from the FT-IR spectroscopy, addition of a modifier oxides to borate and phosphate networks has differing effects on the glass network and hence on alternation of the glass characteristics. The basic units of pure amorphous borate glasses are trigonal  $BO_3$  groups, where as the basic units of pure amorphous phosphate glasses are  $PO_4$  tetrahedra linked through covalent bridging oxygens. In case of borate network, the addition of a modifier oxide increases the degree of polymerization because the neutral trigonal  $BO_{3/2}$  groups are converted to tetrahedral  $BO_{4/2}^-$  species owing to Lewis acid character of boron oxide. Whereas in the phosphate network, the addition of a modifier oxide has a depolymerizing effect; the extra oxygen atoms introduced by the modifier oxides form negative nonbridging oxygen (NBO) sites, whose charge is compensated by the positive charge of the modifier cations [31].

In short, the density alternation of NBO bonds (that have high electronic oxide ion polarizability,  $\alpha_0^{2-}$ ) will lead to such recorded wavering in the glass refractive index values. The alternation can be more understood by reference to Lorentz-Lorenz equation [32] where:

$$\alpha_m = \frac{3M}{4\pi DN_A} \frac{\left(n^2 - 1\right)}{\left(n^2 + 2\right)} \tag{1}$$

where  $\alpha_m$  is the molar polarizability, M is the molar mass of glass composition, D is the glass density and  $N_A$  is the density of oscillators (Avogadro's number). Furthermore, the electronic oxide ion polarizability  $\alpha_0^{2-}$  is function of the molar polarizabilities and which can be considered by the following equation:

$$\alpha_{\mathcal{O}}^{2-} = (\alpha_m - \sum_i N_i \alpha_i) / N_{\mathcal{O}}$$
 (2)

where  $N_0$  and  $N_i$  are the numbers of oxygen ions and cations, respectively. Here  $\alpha_i$  is the electronic polarizabilities of a cation. The values of  $N_0$  and  $N_i$  can be obtained from chemical compositions. Thus, as the ratio of Li<sub>2</sub>O/CaO in the studied glasses composition changes the electronic polarizability of the oxygen ion changes as well describing the obtained

behavior of the glass refractive index.

Duffy et al. [33, 34] proposed the concept of optical basicity which permits a comparison of the acid-base character of oxides. The optical basicity,  $\Lambda$ , of an oxidic medium, is the average electron donor power of all the oxide atoms comprising the medium. Increasing basicity results in increasing the negative charge on the oxygen atoms and, accordingly, increasing covalency in the oxygen-cation bonding. The oxygen (oxide ion) polarizability expresses the degree of negative charge on oxygen atoms while optical basicity is a measurement of this charge [8, 35]. The intrinsic relationship between the optical basicity of the oxide medium and the electronic polarizability of the oxide ion as follows [36]:

$$\Lambda = 1.67 \left( 1 - \frac{1}{\alpha_0^{2-}} \right) \tag{3}$$

An increase in the oxide ion polarizability (due to NBO that impinge on the cation species) leads to an increase in density of the negative charges on the oxygen atoms in the studied glass system and accordingly an increase in glass index.

### **Positron Annihilation Spectroscopy**

In positron annihilation lifetime, the best fit is obtained for three-component analysis of the PAL spectra. The value of positron lifetime  $\tau$  distinguishes the different types of defects. The nonbridging oxygen (NBO) associated with metal cations is considered to have a high electron density than the bridging oxygen [37, 38]. A value in the range of 200–400 ps belongs to bulk and dislocations such as NBOs [39-41]. NBO may be represented as a center of negative charge with a high electron density [37]. The long-lived component  $\tau_3$  is due to the pick-off annihilation of the o-Ps in the free-volume holes of different sizes. On the other hand, the intensity,  $I_3$ , is proportional to the number of these free-volume holes.

According to Tao–Eldrup model [42], knowledge of the o-Ps lifetime,  $\tau_3$ , allows obtaining the average radius, R, of the nano-holes, in spherical approximation as:

$$\tau_3 = 0.5 \left[ 1 - \frac{R}{R + \Delta R} + \frac{1}{2\pi} \sin(\frac{2\pi R}{R + \Delta R}) \right]^{-1}$$
 (4)

 $au_3$  in the range of nanoseconds is related to the existence of free volume, and  $\Delta R = 0.1656$  nm is the thickness of the homogeneous electron layer in which the positron annihilates [43]. Also, the average size of the o-Ps hole volume  $V_f = 4\pi R^3/3$ .

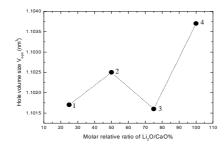


Figure 3: Variation of Free Volume against the Percentage of Molar Ratio of Li<sub>2</sub>O/CaO

Figure 3 shows the effect of the relative percentage of the ratio of  $\text{Li}_2\text{O}/\text{CaO}$  on the free volume size (Vo-ps) measurements of the prepared samples. Similar to what have been seen from fluctuation in the glass indices, a fluctuation in the free volume size is recorded as ratio of  $\text{Li}_2\text{O}/\text{CaO}$  increases. In comparison to sample 1 where the ratio of  $\text{Li}_2\text{O}/\text{CaO}$  = 25 %, an increase in the free volume size is observed at  $\text{Li}_2\text{O}/\text{CaO} = 50$  % (sample 2). The recorded increase in free volume size may be attributed to the location of positrons in the surrounding of an increased number of NBO bonds. The addition of  $\text{Li}_2\text{O}/\text{CaO} = 75$  % (sample 3) causes a significant decrease in the free volume size with lower intensity  $I_2$ . The recorded increase in the intensity for sample 4 indicates the increase in the concentration of NBOs with the increase in the ratio of  $\text{Li}_2\text{O}/\text{CaO} = 100$ % with zero CaO content.

The formation of one  $MO_4$  tetrahedron requires an additional oxygen atom or two single bonded oxygen ions. It can be provided by an alkali molecule such as  $Li_2O$  due to its ionic property [44]. The boric oxide is built up of  $BO_3$  triangles and upon adding divalent oxides, such as CaO, the coordination number of the boron changes from  $SP^3$  tetrahedral  $BO_4$  to from  $SP^2$  planar  $BO_3$ , preserving the B–O bonding without the creation of non-bridging oxygen ions [8].

So the introduction of divalent oxides (CaO) causes a significant formation of the  $BO_3$  groups with a lower coordination number, pursuing an equilibrium reaction  $BO_{3/2} \leftrightarrow BO_{4/2} + NBO$ . When the divalent ions are present in amounts greater than the number of the interstices available, they will act partly as bridges between the adjacent  $BO_4$  tetrahedral. The CaO ions will thus become enclosed in the structural interstices or will act as bridges between the network forming units.

Furthermore, the average coordination number and molar volume are strongly correlated [45]. Therefore, it is expected that, with the insertion of Li<sub>2</sub>O instead of CaO a transformation of a part of BO<sub>3</sub> groups into BO<sub>4</sub> structural unit will take place as it will be seen through the FT-IR measurements. It produces an assured increase in the number of nonbridging oxygen bonds (NBO) relative to the bridging bonds (BO). According to Sun [46], the single bond strength of B–O bond is 498 kJ/mol for BO<sub>3</sub> structural unit and 373 kJ/mol for BO<sub>4</sub> structural unit. It also means the replacement of strong oxygen linkages triagonal units BO<sub>3/2</sub> by weaker anionic tetrahedral unit linkages BO<sup>-</sup>4/2. Thus, the fluctuation in NBO relative density manifests itself clearly by the periodic variation in the glass molar volume and hence the calculated free volume size.

#### FT-IR Structure Analysis

Fourier transform infrared spectroscopy (FTIR) is a technique for collecting and analyzing the obtained IR spectra. The role of the mixed alkali effect on the NBO<sub>s</sub>/BO<sub>s</sub> ratio is found out by analysis of FTIR spectroscopy and positron annihilation spectroscopy. It is also noted the influence of these modifier ions on the structural disorder of the glass matrix. Boron oxygen network can be in the form of planar BO<sub>3</sub> and/or tetrahedral BO<sub>4</sub>.

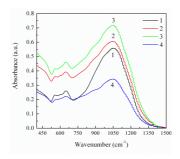


Figure 4: FT-IR Spectra of the different Investigated Glass Compositions

As shown in Fig. 4, the absorption bands are assigned for different stretching and bending vibrations like specific vibrations of B-O-B bending vibrations, O<sub>3</sub>B-O-BO<sub>4</sub> bending vibrations, B-O stretching vibration of BO<sub>4</sub> units in tri-, tetra- and penta- borate groups, B-O stretching vibration of trigonal BO<sub>3</sub> units in boroxol rings, B-O stretching vibrations of BO<sub>3</sub> units in metaborate, pyroborate and orthoborate groups and specific vibrations of B-O-B bending vibrations. The role of the modifier ions (Li<sub>2</sub>O and CaO) on the BO<sub>4</sub>/BO<sub>3</sub> ratio is found out by FTIR analysis. It is also noted the influence of these modifier ions on the structural disorder of the glass matrix. The results of Infrared absorption spectra of the investigated borophosphate glasses in the range extended from 1400 to 400 cm<sup>-1</sup> can be discussed as follow:

- Boron has three vibrational absorption bands at 1200-1600, 800-1200 and at.700 cm<sup>-1</sup>.
- Planar BO<sub>3</sub> gives four fundamental bands around 950, 750, 1250 and 600 cm<sup>-1</sup>.
- Tetrahedral BO<sub>4</sub> unit also gives four bands around 1000, 900, 600, 550 cm<sup>-1</sup>.
- The broad and strong band appeared at 1020-1030 cm<sup>-1</sup> is due to BO<sub>4</sub> tetrahedra.
- The strong band at 700 cm<sup>-1</sup> is due to the bending vibration of B-O-B linkages of BO<sub>3</sub> units [47, 48].
- The weak band observed at 520 cm<sup>-1</sup> is due to vibration of BO<sub>4</sub> tetrahedra.
- It could be concluded that the introduction of Li<sub>2</sub>O transforms some of the BO<sub>3</sub> triangles to BO<sub>4</sub> tetrahedra with the formation of non-bridging oxygen atoms.
- Furthermore, infrared absorption bands of vibration mode P=O superposed with (PO<sub>2</sub>)<sub>as</sub> mode are observed with band assigned to (P-O<sup>-</sup>)<sub>as</sub> vibration in Q<sup>2</sup> (1097 cm<sup>-1</sup>) [49, 50].
- It is also attributed to vibration of Q<sup>0</sup> tetrahedra at 1017 cm<sup>-1</sup> [51-53] and to deformation mode of P-O<sup>-</sup> groups (530 cm<sup>-1</sup>) [54], respectively.
- With increasing Li<sub>2</sub>O on the expense of CaO, the appearance of the band at nearly 913-937 cm<sup>-1</sup> is an indication of the presence of (P-O-P) as bridging oxygen bonds (BOs).
- Appearance of the band at 1045 cm<sup>-1</sup> with increasing Li<sub>2</sub>O proves that the Q<sup>0</sup> tetrahedron exists.
- Disappearances of the band at 990 cm<sup>-1</sup> indicates that Q<sup>1</sup> structure units are decreased.
- So the structure of phosphate network approaches to metaphosphates with mainly Q<sup>2</sup> structural unit.
- When  $\text{Li}_2\text{O}$  is present in larger quantities the  $\pi$ -bond of P=O may be ruptured with the creation of non-bridging oxygen ions.
- Addition of an alkali oxide shifts the reaction  $Q^2 \rightarrow 2Q^1 + Q^0$  to the right direction. Substitution of Li<sub>2</sub>O with CaO depolymerizes the phosphate glass network by systematic conversion of  $Q^2$  structural units into  $Q^1$  structural units [55].
- Conversion Q<sup>2</sup> to Q<sup>1</sup> is taking place due to breaking P-O-P linkages.
- For further addition of Li<sub>2</sub>O, Q<sup>1</sup> structural units get converted into Q<sup>0</sup> structural units and glass network becomes more depolymerized.

• Addition of Li<sub>2</sub>O to phosphate glass in the expense of CaO can decrease the cross-links between phosphate chains [56]. As alkali ion Li<sub>2</sub>O increases it can lead to the breaking of P-O-P linkages and the creation of non-bridging oxygen atoms in the glass [57].

- Increasing the creation of non-bridging oxygens increases the polarization and hence the glass reflectance, R.
- The spectra of indicate that there is no evidence for the presence of characteristic bands of Fe<sub>2</sub>O<sub>3</sub>.
- According to a proposed model [58] the vibrations for Fe<sup>2+</sup> ions should be at 230 cm<sup>-1</sup> and for Fe<sup>3+</sup> ions should be at 290 cm<sup>-1</sup>.
- Furthermore, Mössbauer measurements [59] suggested that, Fe<sup>2+</sup> and Fe<sup>3+</sup> ions occupy sites with octahedral or distorted octahedral coordination without participation in the host glass network giving the reason for absent of absorption bands of Fe–O vibrations.
- It concludes that, vibration of Fe–O may appear below 300 cm<sup>-1</sup> or the band for the Fe–O vibrations may have superimposed or masked with the vibration of O=P–O linkage [60].

### **CONCLUSIONS**

Mixed alkali effect has proven its ability to change the structure and optical properties of alkali borophosphate glasses containing Fe<sub>2</sub>O<sub>3</sub>. The increase in molar percentage of the ratio Li<sub>2</sub>O/CaO induced a periodic change in the ratio of NBO/BO. Such change in the NBO/BO stimulates a similar change in the glass refractive indices and available free volume. The considerable molar and oxide ion polarizabilities of NBO are the motivation for the obtained increase in refractive index value. Furthermore, adjusting the ratio of NBO/BO through the alteration in Li<sub>2</sub>O/CaO leads to a variation in the glass free volume size associated with the change in the glass molar volume.

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